

ROBIN NULLITY IN MODE $|k| = 1$ AND ASYMPTOTIC RADIUS OF THE CRITICAL SPHERICAL CATENOID

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ABSTRACT. For each parameter $a > 1/2$, the critical spherical catenoid Σ_a is a rotationally symmetric, free boundary minimal annulus in a geodesic ball $B^3(r(a)) \subset \mathbb{H}^3$, in the family introduced by Mori [9] and reconsidered by do Carmo–Dajczer [3] and recently by Medvedev [8]. In this short note we establish three analytic results about Σ_a :

(I) *Robin nullity and index in mode $|k| = 1$.* The Robin nullity of the Jacobi operator $L_{\Sigma_a} = \Delta_g + (|\text{II}|^2 - 2)$ in angular Fourier mode $|k| = 1$ equals 2 for every $a > 1/2$, with kernel spanned by the Killing–Jacobi fields associated to the rotations $L_{12}, L_{13} \in \mathfrak{so}(3, 1)$ that fix the geodesic axis of Σ_a and send $\partial B^3(r(a))$ to itself. The radial profile of these Jacobi fields admits the closed form

$$f_*(s) = \partial_s \Phi_a^0(s, 0) = \frac{d}{ds} [A(s) \cosh \varphi(s)] = \sinh r(s) \cdot r'(s),$$

where $r(s) = \text{dist}_{\mathbb{H}^3}(p_0, \Phi_a(s, 0))$. As a consequence of Sturm–Liouville theory and the structure of the zeros of f_* , the Robin Morse index of Σ_a in mode $|k| = 1$ equals 2 for every $a > 1/2$, refining from the analytic side the lower bound $\text{ind}(\Sigma_a) \geq 4$ of Medvedev [8].

(II) *Asymptotic radius.* The boundary radius admits the asymptotic expansion

$$r(a) = \frac{3}{2} \log a + d_\infty + o(1) \quad (a \rightarrow \infty), \quad d_\infty = \log \frac{\sqrt{2} \Gamma(1/4)^2}{\pi^{3/2}} = \log \frac{2\sqrt{2}\pi}{\Gamma(3/4)^2}.$$

The closed form for d_∞ follows from a closed evaluation of the improper integral $I_\infty = \int_0^\infty \cosh(2t)^{-3/2} dt$ via the Beta function.

(III) *Degenerate limit.* As $a \rightarrow (1/2)^+$, $r(a) = c_* \sqrt{a - 1/2} (1 + o(1))$ with $c_* = \sigma_* \cosh \sigma_*$, where σ_* is the unique positive fixed point of $\sigma = \coth \sigma$.

The proof of (I) follows the mode-by-mode strategy of Devyver [2] for the Euclidean critical catenoid, with the group $\mathfrak{so}(3, 1)$ replacing $\mathfrak{so}(3)$, supplemented by the closed-form identification $f_* = \partial_s \Phi^0$ specific to the hyperbolic ambient. The proof of (II) is a Laplace-type asymptotic analysis of the implicit free boundary condition.

Keywords: Free boundary minimal surfaces; Spherical catenoid; Jacobi operator; Robin nullity; Killing fields; Asymptotic geometry; Gamma function values.

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1. INTRODUCTION

Let \mathbb{H}^3 denote three-dimensional hyperbolic space, modelled as the upper sheet of the unit hyperboloid in Minkowski 4-space:

$$\mathbb{H}^3 = \{x \in \mathbb{R}^{3,1} : \langle x, x \rangle_L = -1, x_0 > 0\}, \quad \langle x, y \rangle_L = -x_0 y_0 + \sum_{i=1}^3 x_i y_i,$$

and let $B^3(r) \subset \mathbb{H}^3$ be the closed geodesic ball of radius $r > 0$ centred at $p_0 = (1, 0, 0, 0)$. A *free boundary minimal surface* (FBMS) $\Sigma \subset B^3(r)$ is a properly immersed surface with zero mean curvature in \mathbb{H}^3 meeting $\partial B^3(r)$ orthogonally along $\partial \Sigma$.

The study of free boundary minimal surfaces (FBMS) in geodesic balls has been driven, in the Euclidean case, by the discovery of Fraser and Schoen [4] that the critical catenoid arises as an extremal surface for the first normalized Steklov eigenvalue, establishing a deep link between FBMS and spectral optimization; the uniqueness of the critical catenoid among embedded free boundary minimal annuli in the Euclidean ball is the content of the Fraser–Li conjecture [5]. A central invariant in this circle of ideas is the Morse index: for the Euclidean critical catenoid it

was computed to equal 4, independently and by different methods, by Devyver [2], Smith–Zhou [10], and Tran [12]. The extension of this framework to the space forms \mathbb{S}_+^n and \mathbb{H}^n was initiated by Lima–Menezes [7], who replaced the Steklov problem by a Robin-type eigenvalue problem, and developed by Medvedev [8], who proved the index lower bound $\text{ind}(\Sigma_a) \geq 4$ for the critical spherical catenoid in \mathbb{H}^3 and conjectured equality [8, Remark 5.6]; general upper bounds for the index of FBMS in terms of spectral data were obtained by Lima [6]. The present note contributes three analytic results that sharpen this picture for Σ_a : an exact modal computation of the Robin nullity and index in mode $|k| = 1$, refining Medvedev’s lower bound from the analytic side, together with the precise asymptotic and degenerate geometry of the boundary radius $r(a)$.

1.1. The critical spherical catenoid family. For each $a > 1/2$ there is a one-parameter family of rotationally symmetric FBMS-annuli $\Sigma_a \subset B^3(r(a)) \subset \mathbb{H}^3$, and such family is classically known (Mori [9], do Carmo–Dajczer [3]); it has been revisited by Lima–Menezes [7] and Medvedev [8], who shows that Σ_a has Morse index at least 4 and conjectures equality. Non-rotational FBMS-annuli in geodesic balls of \mathbb{H}^3 have recently been constructed via bifurcation by Cerezo [1]¹. The family $\{\Sigma_a\}_{a>1/2}$ interpolates between the “infinitely thin” degenerate configuration as $a \rightarrow (1/2)^+$ (where the meridian section collapses) and the “unbounded” configuration as $a \rightarrow \infty$ (where the boundary radius diverges).

The parametrization, in the form given by Medvedev [8], reads

$$(1) \quad \Phi_a(s, \theta) = (A(s) \cosh \varphi(s), A(s) \sinh \varphi(s), B(s) \cos \theta, B(s) \sin \theta), \quad (s, \theta) \in \mathbb{R} \times S^1,$$

with

$$(2) \quad A(s)^2 = a \cosh(2s) + \frac{1}{2}, \quad B(s)^2 = a \cosh(2s) - \frac{1}{2},$$

and angular profile

$$(3) \quad \varphi(s) = \int_0^s \frac{\sqrt{a^2 - 1/4}}{A(t)^2 B(t)} dt.$$

A direct computation (Section 2) shows that $\Phi_a : [-s_0, s_0] \times S^1 \rightarrow \mathbb{H}^3$ is an immersion with mean curvature identically zero, induced metric $g = ds^2 + B(s)^2 d\theta^2$, and squared norm of the second fundamental form

$$(4) \quad |\text{II}|^2(s) = 2 \left(\frac{B''(s)}{B(s)} - 1 \right).$$

The free boundary condition selects a discrete value $s_0 = s_0(a) > 0$ and defines the boundary radius $r(a) := \text{arccosh}(A(s_0) \cosh \varphi(s_0))$.

1.2. Statement of results.

Proposition 1.1 (Robin nullity in mode $|k| = 1$). *For every $a > 1/2$, the eigenvalue $\lambda = 0$ of the Robin Jacobi problem*

$$L_{\Sigma_a} u = \lambda u \text{ in } \Sigma_a, \quad \partial_\eta u = \coth(r(a)) u \text{ on } \partial \Sigma_a,$$

restricted to functions in the angular Fourier mode $|k| = 1$ has multiplicity exactly 2. The 2-dimensional kernel is spanned by the Killing–Jacobi fields

$$\langle L_{12}, \nu \rangle_L = f_*(s) \cos \theta, \quad \langle L_{13}, \nu \rangle_L = f_*(s) \sin \theta,$$

where $L_{12}, L_{13} \in \mathfrak{so}(3, 1)$ generate the rotations of $\mathbb{R}^{3,1}$ in the (x_1, x_2) - and (x_1, x_3) -planes, respectively, and f_ is given explicitly in (21).*

¹The rotational surfaces called *free boundary hyperbolic catenoids* by Cerezo [1] are exactly the critical spherical catenoids Σ_a studied here: in the classification of do Carmo–Dajczer [3] they are the rotation hypersurfaces of *spherical* type ($\delta = +1$), whose parallels are geodesic spheres. There the qualifier “hyperbolic” designates the ambient space \mathbb{H}^3 , not the do Carmo–Dajczer type (for which “hyperbolic” is the distinct $\delta = -1$ case).

Theorem 1.2 (Asymptotic radius). *The map $a \mapsto r(a)$ is real-analytic on $(1/2, \infty)$, and as $a \rightarrow \infty$ it satisfies*

$$(5) \quad r(a) = \frac{3}{2} \log a + d_\infty + o(1),$$

where the constant d_∞ is given by the closed form

$$(6) \quad d_\infty = \log \frac{\sqrt{2} \Gamma(1/4)^2}{\pi^{3/2}} = \log \frac{2\sqrt{2\pi}}{\Gamma(3/4)^2}.$$

Proposition 1.3 (Degenerate limit). *As $a \rightarrow (1/2)^+$,*

$$(7) \quad r(a) = c_* \sqrt{a - 1/2} (1 + o(1)), \quad c_* := \rho_* + \frac{1}{\rho_*},$$

where ρ_* is the unique positive root of the transcendental equation

$$(8) \quad \operatorname{arcsinh}(\rho) = \frac{\sqrt{1 + \rho^2}}{\rho}.$$

In particular $r(a) \rightarrow 0$ at rate $r(a) = O(\sqrt{a - 1/2})$.

Remark 1.4. The transcendental equation (8) simplifies under the substitution $\sigma := \operatorname{arcsinh} \rho$, equivalently $\rho = \sinh \sigma$ and $\sqrt{1 + \rho^2} = \cosh \sigma$: (8) becomes

$$(9) \quad \sigma = \coth \sigma,$$

the (unique) positive fixed point of the hyperbolic cotangent. In these variables the leading constant simplifies further: using $\rho_*^2 + 1 = \cosh^2 \sigma_*$ and $\sigma_* \sinh \sigma_* = \cosh \sigma_*$ from (9),

$$(10) \quad c_* = \rho_* + \rho_*^{-1} = \frac{\cosh^2 \sigma_*}{\sinh \sigma_*} = \sigma_* \cosh \sigma_*.$$

The map $\sigma \mapsto \sigma - \coth \sigma$ is real-analytic and strictly increasing on $(0, \infty)$ with limits $-\infty$ and $+\infty$ at the endpoints, giving uniqueness of σ_* (and hence of $\rho_* = \sinh \sigma_*$).

The two equivalent forms of d_∞ in (6) are related by the reflection identity $\Gamma(1/4) \Gamma(3/4) = \pi\sqrt{2}$.

Remark 1.5. Proposition 1.1 extends to the hyperbolic ambient the mode-by-mode analysis of Devyver [2, §3.3], originally carried out for the Euclidean critical catenoid using the $\mathfrak{so}(3)$ -Killing fields. In the hyperbolic setting, only those elements of $\mathfrak{so}(3, 1)$ that preserve the bounding geodesic ball $B^3(r(a))$ produce admissible Jacobi fields; this restricts the relevant Killing subalgebra from a 6-dimensional one to a 3-dimensional one, and the precise computation of which Killing fields contribute to which angular modes occupies Section 4.

Remark 1.6. The closed form (6) is the only step of the proof of Theorem 1.2 that is non-elementary in the sense of relying on the Beta function evaluation $\int_0^\infty \cosh(2t)^{-3/2} dt = \Gamma(3/4)^2 / \sqrt{2\pi}$. The remainder $o(1)$ in (5) can be improved to $O(a^{-1})$ at the expense of a slightly longer asymptotic expansion; we record only the leading-order constant.

1.3. Organization. Section 2 develops the geometric setup and verifies the minimality of Σ_a . Section 3 introduces the relevant Killing–Jacobi fields. Section 4 contains the proof of Proposition 1.1. Section 5 is devoted to the asymptotic analysis and the proof of Theorem 1.2; Proposition 1.3 is established along the way.

2. GEOMETRIC SETUP

In this section we establish the basic geometric properties of Σ_a used throughout the paper.

2.1. Coordinate identities. Throughout we use the parametrization (1)–(3) with $a > 1/2$. Denote $K(a) := \sqrt{a^2 - 1/4}$; this is the conserved quantity associated with the angular momentum of the rotational generator.

Lemma 2.1 (Induced metric). *The pull-back $g = \Phi_a^* g_{\mathbb{H}^3}$ is the diagonal metric*

$$(11) \quad g = ds^2 + B(s)^2 d\theta^2$$

on $\mathbb{R} \times S^1$.

Proof. A direct computation, repeatedly using $A^2 - B^2 = 1$ (which follows from (2)), $\varphi'(s) = K/(A^2 B)$, and the Minkowski inner product:

$$\langle \partial_s \Phi, \partial_s \Phi \rangle_L = -(\partial_s \Phi^0)^2 + (\partial_s \Phi^1)^2 + (\partial_s \Phi^2)^2 + (\partial_s \Phi^3)^2.$$

Using $\partial_s \Phi^0 = A' \cosh \varphi + A \sinh \varphi \cdot \varphi'$ and $\partial_s \Phi^1 = A' \sinh \varphi + A \cosh \varphi \cdot \varphi'$, one obtains

$$-(\partial_s \Phi^0)^2 + (\partial_s \Phi^1)^2 = -(A')^2 + A^2(\varphi')^2.$$

Together with $(\partial_s \Phi^2)^2 + (\partial_s \Phi^3)^2 = (B')^2$, the total reduces, after substituting $A' = a \sinh(2s)/A$, $B' = a \sinh(2s)/B$ and $\varphi' = K/(A^2 B)$, to

$$g_{ss} = -(A')^2 + A^2(\varphi')^2 + (B')^2 = 1,$$

where the last equality is a direct algebraic verification using $A^2 - B^2 = 1$ and $K^2 = a^2 - 1/4$. The off-diagonal term vanishes by symmetry, and $g_{\theta\theta} = B(s)^2$ by inspection. \square

Lemma 2.2 (Minimality and second fundamental form). *The immersion Φ_a has identically vanishing mean curvature, and*

$$(12) \quad |\Pi|^2(s) = 2 \left(\frac{B''(s)}{B(s)} - 1 \right).$$

Proof. Minimality of Σ_a is a classical fact (Mori [9], do Carmo–Dajczer [3], Medvedev [8]); we take it as known.

For (4): let ν be a unit normal to Σ_a in \mathbb{H}^3 . Set $\kappa_s = \Pi(\partial_s, \partial_s)/g_{ss}$ and $\kappa_\theta = \Pi(\partial_\theta, \partial_\theta)/g_{\theta\theta}$, the principal curvatures. By the Gauss equation in \mathbb{H}^3 (sectional curvature -1):

$$K_\Sigma = -1 + \kappa_s \kappa_\theta.$$

On the other hand, for the diagonal metric $g = ds^2 + B^2 d\theta^2$ the intrinsic Gauss curvature is the standard expression $K_\Sigma = -B''/B$. Hence $\kappa_s \kappa_\theta = 1 - B''/B$. By minimality $\kappa_s + \kappa_\theta = 0$, i.e. $\kappa_s = -\kappa_\theta$, so $\kappa_\theta^2 = B''/B - 1$, and $|\Pi|^2 = \kappa_s^2 + \kappa_\theta^2 = 2(B''/B - 1)$, giving (4). \square

Lemma 2.3 (Conserved quantity identity). *On Σ_a ,*

$$(13) \quad |\Pi|^2(s) B(s)^4 = 2K^2,$$

i.e. $|\Pi|^2(s) = 2K^2/B(s)^4$.

Proof. From $B^2 = a \cosh(2s) - 1/2$, differentiating once gives $BB' = a \sinh(2s)$, and differentiating again, $(B')^2 + BB'' = 2a \cosh(2s) = 2B^2 + 1$. Hence

$$(14) \quad BB'' = 2B^2 + 1 - (B')^2.$$

Now $(B')^2 = (BB')^2/B^2 = a^2 \sinh^2(2s)/B^2$, and

$$B^4(B''/B - 1) = B^4 + B^2 - (B')^2 \cdot B^2 = B^2(B^2 + 1) - a^2 \sinh^2(2s).$$

Since $B^2 + 1 = a \cosh(2s) + 1/2$ and $B^2 = a \cosh(2s) - 1/2$,

$$B^2(B^2 + 1) = (a \cosh(2s))^2 - 1/4 = a^2 \cosh^2(2s) - 1/4,$$

hence

$$B^4(B''/B - 1) = a^2 \cosh^2(2s) - 1/4 - a^2 \sinh^2(2s) = a^2 - 1/4 = K^2,$$

where we used $\cosh^2 - \sinh^2 = 1$. Combining with (4) yields (13). \square

Lemma 2.4 (Coordinate equations and free boundary condition). *Each coordinate function $\Phi_a^A : \Sigma_a \rightarrow \mathbb{R}$ satisfies, with the geometric sign convention $\Delta_g = \operatorname{div}_g \nabla$,*

$$(15) \quad \Delta_g \Phi_a^A = 2\Phi_a^A \quad (A = 0, 1, 2, 3).$$

The free boundary condition $\Sigma_a \perp \partial B^3(r(a))$ is equivalent to

$$(16) \quad \nu^0(s_0, \theta) = 0,$$

where ν^0 is the 0-th Minkowski component of the unit normal, and where $s_0 = s_0(a)$ is the unique positive root of

$$(17) \quad \tanh \varphi(s) = \frac{B(s) K(a)}{a \sinh(2s)}.$$

The boundary radius is then $r(a) = \operatorname{arccosh}(A(s_0) \cosh \varphi(s_0))$.

Proof. For (15): in \mathbb{H}^3 , the coordinate functions Φ^A ($A = 0, 1, 2, 3$) satisfy $\Delta_{\mathbb{H}^3} \Phi^A = 3\Phi^A$ (direct computation in geodesic coordinates, or alternatively the identity $\Delta_{\mathbb{H}^n} \Phi^A = n\Phi^A$ for the ambient coordinates of the Lorentz model). Since Σ_a is minimal and of codimension one, the standard relation $\Delta_g f = \Delta_{\mathbb{H}^3} f|_{\Sigma_a} - \operatorname{Hess}_{\mathbb{H}^3}(f)(\nu, \nu)$ holds for any function f on \mathbb{H}^3 . The Hessian of Φ^A in \mathbb{H}^3 satisfies $\operatorname{Hess}_{\mathbb{H}^3}(\Phi^A) = \Phi^A g_{\mathbb{H}^3}$ (a direct computation), so $\operatorname{Hess}_{\mathbb{H}^3}(\Phi^A)(\nu, \nu) = \Phi^A$, and (15) follows.

For the free boundary condition: $\partial B^3(r) \subset \mathbb{H}^3$ is the level set $\{x \in \mathbb{H}^3 : x^0 = \cosh(r)\}$. Its inward \mathbb{H}^3 -unit normal at a point $\Phi_a(s, \theta)$ on the boundary is

$$\eta_{\partial B^3(r)} = \frac{1}{\sinh(r)} (-\partial_0 + \Phi_a^0 \Phi_a),$$

where $\partial_0 = (1, 0, 0, 0)$. The orthogonality condition reads $\langle \nu, \eta_{\partial B^3(r)} \rangle_L = 0$, i.e. $-\nu^0 \cdot (-1) + \Phi_a^0 \langle \nu, \Phi_a \rangle_L = 0$. Since $\langle \nu, \Phi_a \rangle_L = 0$ along Σ_a (being normal in the ambient sense), this reduces to $\nu^0 = 0$, which is (16).

The explicit form (17) follows by direct computation of ν from the cross-product formula in $\mathbb{R}^{3,1}$ and the identity $A^2 - B^2 = 1$. For existence and uniqueness of $s_0 \in (0, \infty)$ for each $a > 1/2$: the left-hand side of (17), $\tanh \varphi(s)$, is strictly increasing from $\tanh \varphi(0) = 0$ towards $\tanh \varphi(\infty) < 1$ as $s \rightarrow \infty$ (since φ is itself strictly increasing and bounded above; the bound follows from the convergence of the integral defining φ). The right-hand side of (17),

$$R(s) := \frac{B(s) K(a)}{a \sinh(2s)},$$

behaves as $R(s) \sim B(0)K(a)/(2as) \rightarrow +\infty$ as $s \rightarrow 0^+$, and $R(s) \rightarrow 0$ as $s \rightarrow \infty$ (since $B(s) \sim a^{1/2} 2^{-1/2} e^s$ and $\sinh(2s) \sim e^{2s}/2$, so $R(s) \sim \operatorname{const} \cdot e^{-s} \rightarrow 0$). The function R is strictly decreasing on $(0, \infty)$: indeed, using $B'/B = a \sinh(2s)/B^2$,

$$\frac{R'(s)}{R(s)} = \frac{B'(s)}{B(s)} - 2 \coth(2s) = \frac{a \sinh(2s)}{B^2} - \frac{2 \cosh(2s)}{\sinh(2s)},$$

and putting the two terms over the common denominator $B^2 \sinh(2s) > 0$ together with $B^2 = a \cosh(2s) - \frac{1}{2}$ and $\sinh^2(2s) = \cosh^2(2s) - 1$, the numerator becomes

$$a \sinh^2(2s) - 2 \cosh(2s) B^2 = -(a \cosh^2(2s) - \cosh(2s) + a).$$

The polynomial $\xi \mapsto a\xi^2 - \xi + a$ has discriminant $1 - 4a^2 < 0$ for $a > 1/2$ and positive leading coefficient, hence is strictly positive for every real ξ . Therefore $R'(s)/R(s) < 0$ on $(0, \infty)$. Consequently the equation $\tanh \varphi(s) = R(s)$ admits a unique positive solution $s_0 = s_0(a)$. \square

3. KILLING–JACOBI FIELDS AND THE VARIATION PRINCIPLE

The Lie algebra of isometries of \mathbb{H}^3 is $\mathfrak{so}(3, 1)$, of dimension 6. Its elements correspond to Killing vector fields on \mathbb{H}^3 . Among these, only the 3-dimensional subalgebra of *spatial rotations* $\mathfrak{so}(3) \subset \mathfrak{so}(3, 1)$, generated by $L_{ij} = x_i \partial_j - x_j \partial_i$ for $1 \leq i < j \leq 3$, fixes the geodesic ball $B^3(r)$ centred at $p_0 = (1, 0, 0, 0)$. The remaining Lorentzian boosts $L_{0i} = x_0 \partial_i + x_i \partial_0$ ($i = 1, 2, 3$) do not preserve $B^3(r)$ and therefore do not generate admissible variations.

Lemma 3.1 (Killing–Jacobi fields). *Let K be any element of $\mathfrak{so}(3) \subset \mathfrak{so}(3, 1)$, viewed as a vector field on \mathbb{H}^3 . The function*

$$u_K := \langle K, \nu \rangle_L : \Sigma_a \rightarrow \mathbb{R}$$

satisfies the Jacobi equation

$$L_{\Sigma_a} u_K = \Delta_g u_K + (|\Pi|^2 - 2)u_K = 0,$$

together with the Robin boundary condition

$$\partial_\eta u_K = \coth(r(a)) u_K \quad \text{on } \partial \Sigma_a.$$

Proof. This is a direct adaptation to \mathbb{H}^3 of the standard variation principle for Jacobi fields generated by ambient Killing fields (cf. [2, §3.3] for the Euclidean case): the flow of K is an isometry of \mathbb{H}^3 that preserves $B^3(r(a))$, and therefore generates a one-parameter family of FBMS through Σ_a . The normal component of the variation vector field is exactly $u_K \nu$, and the Jacobi equation together with the Robin boundary condition is the linearized statement of “minimality is preserved at first order” and “free boundary condition is preserved at first order”. The constant $\coth(r(a))$ in the Robin condition arises from the second fundamental form of $\partial B^3(r(a))$ in \mathbb{H}^3 : geodesic spheres $\partial B^3(r)$ in \mathbb{H}^3 have second fundamental form $\coth(r) g_{\partial B^3(r)}$ (umbilical with constant principal curvature $\coth r$). We omit the detailed verification, which is standard. \square

We focus in particular on L_{12} and L_{13} . The unit normal to Σ_a in \mathbb{H}^3 admits a closed form along the meridian $\theta = 0$, which we now derive.

Lemma 3.2 (Unit normal along $\theta = 0$). *For every $s \in (-s_0, s_0)$, the unit normal ν to Σ_a in \mathbb{H}^3 at the point $\Phi_a(s, 0)$ is given (up to a global choice of orientation, fixed once and for all) by*

$$(18) \quad \begin{aligned} \nu^0(s, 0) &= \frac{K \cosh \varphi(s)}{A(s)} - \frac{a \sinh(2s) \sinh \varphi(s)}{A(s)B(s)}, \\ \nu^1(s, 0) &= \frac{K \sinh \varphi(s)}{A(s)} - \frac{a \sinh(2s) \cosh \varphi(s)}{A(s)B(s)}, \\ \nu^2(s, 0) &= \frac{K}{B(s)}, \quad \nu^3(s, 0) = 0. \end{aligned}$$

Proof. At $\theta = 0$, $\Phi_a(s, 0) = (A \cosh \varphi, A \sinh \varphi, B, 0)$, and the two tangent vectors to Σ_a at this point are

$$T_s = (A' \cosh \varphi + A \varphi' \sinh \varphi, A' \sinh \varphi + A \varphi' \cosh \varphi, B', 0), \quad T_\theta = (0, 0, 0, B).$$

The unit normal $\nu(s, 0) \in T_{\Phi_a} \mathbb{H}^3$ satisfies the four conditions

$$(19) \quad \langle \nu, \Phi_a \rangle_L = 0, \quad \langle \nu, T_s \rangle_L = 0, \quad \langle \nu, T_\theta \rangle_L = 0, \quad \langle \nu, \nu \rangle_L = 1.$$

The third condition gives $B \nu^3 = 0$, hence $\nu^3(s, 0) = 0$. The first two are linear in (ν^0, ν^1, ν^2) :

$$(20) \quad \begin{cases} -A \cosh \varphi \nu^0 + A \sinh \varphi \nu^1 + B \nu^2 = 0, \\ -(A' \cosh \varphi + A \varphi' \sinh \varphi) \nu^0 + (A' \sinh \varphi + A \varphi' \cosh \varphi) \nu^1 + B' \nu^2 = 0. \end{cases}$$

Solving (20) for (ν^0, ν^1) in terms of ν^2 by Cramer’s rule, the relevant 2×2 determinant of the (ν^0, ν^1) -block equals

$$-A \cosh \varphi (A' \sinh \varphi + A \varphi' \cosh \varphi) + A \sinh \varphi (A' \cosh \varphi + A \varphi' \sinh \varphi) = -A^2 \varphi' = -K/B,$$

where the last equality uses $\varphi' = K/(A^2B)$. The resulting expressions, after using $AB' - BA' = a \sinh(2s)/(AB)$ (which follows from $A' = a \sinh(2s)/A$, $B' = a \sinh(2s)/B$ and $A^2 - B^2 = 1$) and $AB\varphi' = K/A$, simplify to

$$\nu^0 = \frac{\nu^2}{AK} [BK \cosh \varphi - a \sinh(2s) \sinh \varphi], \quad \nu^1 = \frac{\nu^2}{AK} [BK \sinh \varphi - a \sinh(2s) \cosh \varphi].$$

Substituting these into the unit-norm condition, $-(\nu^0)^2 + (\nu^1)^2 + (\nu^2)^2 = 1$, and using the identity (proved by direct expansion)

$$[BK \sinh \varphi - a \sinh(2s) \cosh \varphi]^2 - [BK \cosh \varphi - a \sinh(2s) \sinh \varphi]^2 = a^2 \sinh^2(2s) - B^2 K^2,$$

together with $K^2 + a^2 \sinh^2(2s) = a^2 \cosh^2(2s) - 1/4 = A^2 B^2$, one obtains $(\nu^2)^2 = K^2/B^2$. Choosing the orientation $\nu^2(s, 0) = K/B(s) > 0$ (which is consistent and does not vanish since $B(s) > 0$ for $a > 1/2$), the formulas (18) follow. \square

Remark 3.3. The free boundary condition $\nu^0(s_0, 0) = 0$ reads, from (18), $BK \cosh \varphi = a \sinh(2s) \sinh \varphi$ at $s = s_0$, i.e. $\tanh \varphi(s_0) = BK/(a \sinh(2s_0))$, recovering (17).

Lemma 3.4. *With the parametrization (1) and the convention $L_{12} = x_1 \partial_2 - x_2 \partial_1$, $L_{13} = x_1 \partial_3 - x_3 \partial_1$:*

$$\langle L_{12}, \nu \rangle_L = f_*(s) \cos \theta, \quad \langle L_{13}, \nu \rangle_L = f_*(s) \sin \theta,$$

where

$$(21) \quad f_*(s) = \frac{K \sinh \varphi(s) + a B(s) \sinh(2s) \cosh \varphi(s)}{A(s)B(s)}.$$

The function f_* is real-analytic and odd in s , vanishes only at $s = 0$, and is strictly positive on $(0, s_0]$ (hence strictly negative on $[-s_0, 0)$).

Proof. Since L_{12} acts on the (x_1, x_2) -plane, its restriction to the point $\Phi_a(s, \theta)$ is the vector field with components $(0, -x_2, x_1, 0) = (0, -B \cos \theta, A \sinh \varphi, 0)$. Therefore

$$\langle L_{12}, \nu \rangle_L = -B \cos \theta \nu^1(s, \theta) + A \sinh \varphi \nu^2(s, \theta).$$

By the rotational symmetry $\Phi_a(s, \theta + \alpha) = R_\alpha \Phi_a(s, \theta)$, where R_α is the rotation by α in the (x_2, x_3) -plane, the unit normal transforms as $\nu(s, \theta + \alpha) = R_\alpha \nu(s, \theta)$. Applied to the values along $\theta = 0$ provided by Lemma 3.2, this gives

$$\nu^0(s, \theta) = \nu^0(s, 0), \quad \nu^1(s, \theta) = \nu^1(s, 0), \quad \nu^2(s, \theta) = \cos \theta \nu^2(s, 0), \quad \nu^3(s, \theta) = \sin \theta \nu^2(s, 0),$$

the simplification of ν^2, ν^3 being possible because $\nu^3(s, 0) = 0$. Substituting,

$$\langle L_{12}, \nu \rangle_L = [A \sinh \varphi \cdot \nu^2(s, 0) - B \cdot \nu^1(s, 0)] \cos \theta.$$

The same computation for L_{13} yields the analogous expression with $\sin \theta$. The bracket equals

$$\frac{K A \sinh \varphi}{B} - B \left[\frac{K \sinh \varphi}{A} - \frac{a \sinh(2s) \cosh \varphi}{AB} \right] = K \sinh \varphi \frac{A^2 - B^2}{AB} + \frac{a \sinh(2s) \cosh \varphi}{A}.$$

Using $A^2 - B^2 = 1$ and reducing to common denominator AB gives (21).

Parity. In $f_*(s)$ the numerator $K \sinh \varphi(s) + a B(s) \sinh(2s) \cosh \varphi(s)$ is the sum of two odd functions of s ($\sinh \varphi$ is odd because φ is the integral from 0 of a positive even function, and $\sinh(2s)$ is odd; $A, B, \cosh \varphi$ are even, and a, K are constants). The denominator $A(s)B(s)$ is even. Hence f_* is odd.

Sign and zeros. For $s \in (0, s_0]$, all factors $K, \sinh \varphi(s), a, B(s), \sinh(2s), \cosh \varphi(s), A(s)$ are strictly positive (this uses $a > 1/2$ and the strict monotonicity of φ), so $f_*(s) > 0$. By oddness, $f_*(s) < 0$ on $[-s_0, 0)$. Finally, $f_*(0) = 0$ by direct evaluation ($\sinh \varphi(0) = 0$ and $\sinh(0) = 0$). Real-analyticity is inherited from A, B, φ . \square

In particular, f_* is a non-zero solution to the radial ODE in mode $|k| = 1$:

$$(22) \quad f_*'' + \frac{B'}{B} f_*' + \left(|\text{II}|^2 - 2 - \frac{1}{B^2} \right) f_* = 0,$$

satisfying $f_*'(s_0) = \coth(r(a)) f_*(s_0)$.

The closed form (21) admits a structural rewriting which will play a central role in Section 4 and in the applications below.

Corollary 3.5 (Geometric form of f_*). *The radial profile f_* of the Killing–Jacobi fields generated by L_{12}, L_{13} coincides with the s -derivative of the time coordinate along the meridian $\theta = 0$:*

$$(23) \quad f_*(s) = \partial_s \Phi_a^0(s, 0) = \frac{d}{ds} [A(s) \cosh \varphi(s)].$$

Equivalently, denoting $r(s) := \text{dist}_{\mathbb{H}^3}(p_0, \Phi_a(s, 0))$,

$$(24) \quad f_*(s) = \sinh r(s) \cdot r'(s) = \frac{d}{ds} \cosh r(s).$$

Proof. Direct expansion of (21) using the identities $K = A^2 B \varphi'$ (definition of φ') and $a \sinh(2s) = AA'$ (from $A' = a \sinh(2s)/A$) gives

$$f_*(s) = \frac{K \sinh \varphi}{AB} + \frac{a \sinh(2s) \cosh \varphi}{A} = A \varphi' \sinh \varphi + A' \cosh \varphi = \partial_s (A \cosh \varphi),$$

which is (23). Equation (24) follows from $\Phi_a^0(s, 0) = \cosh(\text{dist}_{\mathbb{H}^3}(p_0, \Phi_a(s, 0)))$, which is the standard expression for the geodesic distance in the hyperboloid model: for $p_0 = (1, 0, 0, 0)$ and $x \in \mathbb{H}^3$, $\cosh \text{dist}_{\mathbb{H}^3}(p_0, x) = -\langle p_0, x \rangle_L = x^0$. \square

Remark 3.6. The identity (23) reflects a structural feature of the hyperbolic ambient. The function Φ_a^0 on Σ_a satisfies $\Delta_g \Phi_a^0 = 2\Phi_a^0$ (Lemma 2.4), and hence $\partial_s \Phi_a^0$ at $\theta = 0$ satisfies the θ -differentiated equation

$$(\partial_s \Phi_a^0)'' + \frac{B'}{B} (\partial_s \Phi_a^0)' + \left(\frac{B''}{B} - \left(\frac{B'}{B} \right)^2 - 2 \right) (\partial_s \Phi_a^0) = 0,$$

which coincides with the mode- $|k| = 1$ Jacobi ODE (22) exactly when $(B'/B)^2 - B''/B + 2 = 2 - |\text{II}|^2 + 1/B^2$, equivalently $BB'' + (B')^2 = 2B^2 + 1$ — precisely the Mori catenoid identity (14).

Remark 3.7. The identity $\Delta_g \Phi_a^0 = 2\Phi_a^0$ on Σ_a has a structural origin which extends beyond the specific Mori parametrization. On the ambient \mathbb{H}^3 , the four coordinate functions $\Phi^A : \mathbb{H}^3 \rightarrow \mathbb{R}$ ($A = 0, 1, 2, 3$) inherited from the embedding $\mathbb{H}^3 \hookrightarrow \mathbb{R}^{3,1}$ satisfy the tensorial identity

$$(25) \quad \text{Hess}^{\mathbb{H}^3} \Phi^A(X, Y) = \langle X, Y \rangle_{g_{\mathbb{H}^3}} \cdot \Phi^A \quad \text{for all } X, Y \in T\mathbb{H}^3,$$

i.e. the Hessian is conformal to the ambient metric with conformal factor Φ^A itself; the proof is a direct application of the Gauss formula for $\mathbb{H}^3 \subset \mathbb{R}^{3,1}$ together with the relation $\text{II}_{\mathbb{R}^{3,1}}^{\mathbb{H}^3}(X, Y) = \langle X, Y \rangle_g$. Tracing (25) yields $\Delta_{\mathbb{H}^3} \Phi^A = 3\Phi^A$, and restricting to a 2-dimensional minimal $\Sigma \subset \mathbb{H}^3$ gives $\Delta_g \Phi^A = 2\Phi^A$ as in Lemma 2.4.

The positive sign of the eigenvalue is the analytic signature of the constant negative sectional curvature of the ambient: the analogous identities are $\text{Hess}^{\mathbb{R}^n} \Phi^A = 0$ in flat space (so ambient coordinates restrict to harmonic functions on minimal surfaces in \mathbb{R}^n , eigenvalue 0) and $\text{Hess}^{S^n} \Phi^A = -\Phi^A g$ in the round sphere (eigenvalue $-k$ on k -dimensional minimal submanifolds). The Tashiro-type equation (25), with $f = \Phi^A$ playing the role of the unknown, is in fact a classical hallmark of constant-curvature space forms (cf. [11]). From this perspective, Corollary 3.5 expresses the radial Killing–Jacobi profile f_* as the meridian derivative of a curvature-signed eigenfunction of the intrinsic Laplacian: a manifestly hyperbolic phenomenon, with no Euclidean counterpart.

4. PROOF OF PROPOSITION 1.1

We now prove that the Robin kernel of the Jacobi operator $L_{\Sigma_a} = \Delta_g + (|\mathbb{II}|^2 - 2)$ restricted to functions of the form

$$(26) \quad u(s, \theta) = f(s) \cos \theta \quad \text{or} \quad u(s, \theta) = f(s) \sin \theta$$

has dimension exactly 2 for every $a > 1/2$.

Substituting (26) into the Jacobi equation $L_{\Sigma_a} u = 0$ with the Robin boundary condition $\partial_\eta u = \coth(r(a))u$ at $s = \pm s_0$ reduces, by separation of variables in the metric $g = ds^2 + B(s)^2 d\theta^2$, to the radial boundary value problem

$$(27) \quad \begin{cases} f'' + \frac{B'}{B} f' + \left(|\mathbb{II}|^2 - 2 - \frac{1}{B^2} \right) f = 0, & s \in (-s_0, s_0), \\ f'(\pm s_0) = \pm \coth(r(a)) f(\pm s_0), \end{cases}$$

with the sign convention that $f'(-s_0)$ denotes the derivative on the (right) interior side at $s = -s_0$, i.e. the outward normal derivative is $-f'(-s_0)$ at the left endpoint.

The differential operator and the boundary conditions in (27) are invariant under $s \mapsto -s$, so the radial kernel decomposes as

$$(28) \quad \ker^{\text{rad}} = \ker_{\text{odd}}^{\text{rad}} \oplus \ker_{\text{even}}^{\text{rad}}.$$

Moreover, in each parity class, the radial ODE has a 2-dimensional solution space, of which the parity restriction selects a 1-dimensional subspace at the ODE level, and the Robin boundary condition selects a (possibly empty) further restriction. Thus

$$(29) \quad \dim \ker_{\text{odd}}^{\text{rad}} \leq 1, \quad \dim \ker_{\text{even}}^{\text{rad}} \leq 1.$$

The proof of Proposition 1.1 reduces to the two assertions

$$(30) \quad \dim \ker_{\text{odd}}^{\text{rad}} = 1, \quad \dim \ker_{\text{even}}^{\text{rad}} = 0.$$

Lemma 4.1 (Odd radial kernel is one-dimensional). *For every $a > 1/2$, the space $\ker_{\text{odd}}^{\text{rad}}$ has dimension exactly 1, and is spanned by the function f_* from (21).*

Proof. By Lemma 3.4, the function $\langle L_{12}, \nu \rangle_L = f_*(s) \cos \theta$ lies in $\ker(L_{\Sigma_a})$ and satisfies the Robin boundary condition. Its radial profile f_* is odd in s and strictly positive on $(0, s_0]$ (Lemma 3.4), in particular non-zero. Hence $\dim \ker_{\text{odd}}^{\text{rad}} \geq 1$. The opposite inequality is (29). \square

The non-trivial step is to exclude the even-radial sector.

Lemma 4.2 (Even radial kernel is trivial). *For every $a > 1/2$, $\dim \ker_{\text{even}}^{\text{rad}} = 0$.*

Proof. Suppose, for contradiction, that there exists a non-zero $f_e \in \ker_{\text{even}}^{\text{rad}}$. Together with $f_* \in \ker_{\text{odd}}^{\text{rad}}$, the pair (f_e, f_*) is linearly independent, so it forms a fundamental system of solutions of the radial ODE on $(-s_0, s_0)$.

Consider the Wronskian

$$W(s) := f_*(s) f_e'(s) - f_e(s) f_*'(s).$$

The radial ODE in (27) can be written in self-adjoint form as

$$(B(s) f'(s))' + B(s) \left(|\mathbb{II}|^2 - 2 - \frac{1}{B^2} \right) f(s) = 0,$$

which yields the Abel-type identity

$$(31) \quad B(s) W(s) = \text{const} =: C_W.$$

We now evaluate (31) at $s = s_0$. Both f_* and f_e satisfy the Robin condition at s_0 :

$$f_*'(s_0) = \coth(r) f_*(s_0), \quad f_e'(s_0) = \coth(r) f_e(s_0),$$

so

$$W(s_0) = f_*(s_0) (\coth(r) f_e(s_0)) - f_e(s_0) (\coth(r) f_*(s_0)) = 0,$$

hence $C_W = B(s_0) \cdot 0 = 0$, and so $W(s) \equiv 0$ on $(-s_0, s_0)$.

We now derive a contradiction by evaluating W at $s = 0$. Since f_* is odd and non-trivial, $f_*(0) = 0$ and $f'_*(0) \neq 0$ (otherwise $f_* \equiv 0$ by uniqueness for the second-order ODE with zero Cauchy data). Symmetrically, since f_e is even and non-trivial, $f'_e(0) = 0$ and $f_e(0) \neq 0$. Therefore

$$W(0) = f_*(0)f'_e(0) - f_e(0)f'_*(0) = -f_e(0)f'_*(0) \neq 0,$$

contradicting $W \equiv 0$. Hence no non-trivial f_e exists. \square

Proof of Proposition 1.1. Combining Lemmas 4.1 and 4.2, the radial kernel (28) satisfies $\dim \ker^{\text{rad}} = 1$, with $\ker^{\text{rad}} = \mathbb{R}\langle f_* \rangle$. In the full mode $|k| = 1$ (which is the two-dimensional space spanned by $\cos \theta$ and $\sin \theta$ as angular factors), the kernel is therefore

$$\mathbb{R}\langle f_*(s) \cos \theta \rangle \oplus \mathbb{R}\langle f_*(s) \sin \theta \rangle,$$

which has dimension 2. By Lemma 3.4, these are identified with $\langle L_{12}, \nu \rangle_L$ and $\langle L_{13}, \nu \rangle_L$ respectively. \square

Remark 4.3. The explicit identification of the kernel with Killing–Jacobi fields makes the result *geometric*: the 2-dimensional kernel in mode $|k| = 1$ reflects the 2-dimensional family of “small” isometric deformations of Σ_a inside $B^3(r(a))$, namely those that rotate the geodesic ball axis of Σ_a infinitesimally about the x_2 - and x_3 -axes. These do not produce new FBMS; they are infinitesimal congruences.

Corollary 4.4 (Robin Morse index in mode $|k| = 1$). *For every $a > 1/2$, the Robin Morse index of Σ_a in mode $|k| = 1$ equals 2. Equivalently, the radial Sturm–Liouville problem (27) (with $\lambda = 0$ replaced by λ) has exactly one negative eigenvalue $\mu_0(a) < 0$, and the next eigenvalue is $\mu_1(a) = 0$.*

Proof. The radial mode- $|k| = 1$ problem on $(-s_0, s_0)$ is a regular Sturm–Liouville problem in self-adjoint form

$$-(B(s) f'(s))' + B(s) \left(2 + \frac{1}{B(s)^2} - |\text{II}|^2(s) \right) f(s) = \mu B(s) f(s),$$

with separated Robin boundary conditions (27), and with $B(s) > 0$ on the closed interval and continuous coefficients. By classical Sturm–Liouville theory, the spectrum is discrete $\mu_0 < \mu_1 < \mu_2 < \dots$ with $\mu_n \rightarrow +\infty$, and the eigenfunction ϕ_n corresponding to μ_n has exactly n zeros in the open interval $(-s_0, s_0)$.

By Lemma 3.4, f_* is a non-trivial radial Robin eigenfunction with eigenvalue 0, having exactly one zero in $(-s_0, s_0)$ (at $s = 0$). Hence f_* is the second eigenfunction, i.e. $\mu_1(a) = 0$, and $\mu_0(a) < \mu_1(a) = 0$ is strictly negative.

In the full mode $|k| = 1$, each radial eigenvalue μ_n has 2-dimensional angular multiplicity (spanned by $\cos \theta$ and $\sin \theta$). Hence the Robin Jacobi operator restricted to mode $|k| = 1$ has exactly two negative eigenvalues (both equal to $\mu_0(a)$) and a 2-dimensional kernel (recovering Proposition 1.1). \square

Remark 4.5. The notion of *Robin Morse index* used in Corollary 4.4 coincides with the Morse index $\text{ind}(\Sigma_a)$ as defined by Medvedev [8, Definition 5.1]: the maximal dimension of a subspace of $C^\infty(\Sigma_a)$ on which the second-variation quadratic form

$$S(u, u) = \int_{\Sigma_a} (|\nabla u|^2 - (|\text{II}|^2 - 2)u^2) dA - \coth(r(a)) \int_{\partial \Sigma_a} u^2 dL$$

is negative definite. By integration by parts, $S(u, u)$ equals $-\int_{\Sigma_a} u L_{\Sigma_a} u dA$ on functions satisfying the Robin condition $\partial_\eta u = \coth(r(a)) u$ on $\partial \Sigma_a$, so the negative directions of S are exactly the Robin eigenfunctions of L_{Σ_a} with positive eigenvalue; equivalently, the Robin Morse index is the number of negative eigenvalues of the second-variation form S (i.e. of the operator $-L_{\Sigma_a}$ under the Robin boundary condition). Being the dimension of the negative subspace of the geometric form S , the index is independent of the sign convention adopted for Δ_g (Medvedev: $\Delta_g = -\text{div}_g \nabla$; here: $\Delta_g = +\text{div}_g \nabla$): the convention only flips the signs of the eigenvalues of the Jacobi operator, hence whether the index is read off the positive

or the negative part of its spectrum, not the dimension of the negative subspace of S . The mode-by-mode decomposition employed in this paper is the orthogonal Fourier decomposition $L^2(\Sigma_a) = \bigoplus_{k \in \mathbb{Z}} L^2((-s_0, s_0), B(s) ds) \otimes \mathbb{C}\langle e^{ik\theta} \rangle$ on the product domain $(-s_0, s_0) \times S^1$, which is preserved by L_{Σ_a} (rotational invariance) and by the Robin boundary condition.

Remark 4.6. Combined with the lower bound $\text{ind}(\Sigma_a) \geq 4$ established by Medvedev [8], Corollary 4.4 implies that the Robin Morse index of Σ_a in modes $|k| \neq 1$ is at least 2. The conjectural equality $\text{ind}(\Sigma_a) = 4$ would then follow from showing that this lower bound of 2 is also an upper bound; see Section 6.

5. PROOF OF THEOREM 1.2

We analyze the implicit free boundary equation (17) in the regime $a \rightarrow \infty$ to derive the asymptotic expansion (5). We also record the elementary computation establishing Proposition 1.3 in passing.

5.1. **The improper integral.** Recall the angular profile φ from (3):

$$\varphi(s) = \int_0^s \frac{K(a)}{A(t)^2 B(t)} dt, \quad K(a) = \sqrt{a^2 - 1/4}.$$

For a large and $s \rightarrow \infty$ with s allowed to grow,

$$A(t)^2 B(t) = (a \cosh(2t) + \frac{1}{2}) \cdot (a \cosh(2t) - \frac{1}{2})^{1/2} = a^{3/2} \cosh(2t)^{3/2} \cdot (1 + O(a^{-1})),$$

uniformly in $t \geq 0$, and $K(a) = a(1 + O(a^{-2}))$. Hence

$$(32) \quad \varphi(s) = \frac{1}{\sqrt{a}} \int_0^s \cosh(2t)^{-3/2} dt \cdot (1 + O(a^{-1})).$$

Lemma 5.1 (Closed form for the improper integral). *One has*

$$(33) \quad I_\infty := \int_0^\infty \cosh(2t)^{-3/2} dt = \frac{\Gamma(3/4)^2}{\sqrt{2\pi}}.$$

Proof. Substitute $u = e^{-4t}$, so that $e^{2t} = u^{-1/2}$, $dt = -du/(4u)$, and $\cosh(2t) = (u^{-1/2} + u^{1/2})/2 = (1+u)/(2\sqrt{u})$. Then

$$\cosh(2t)^{-3/2} = \frac{2\sqrt{2} u^{3/4}}{(1+u)^{3/2}},$$

and the integral becomes

$$I_\infty = \int_0^1 \frac{2\sqrt{2} u^{3/4}}{(1+u)^{3/2}} \cdot \frac{du}{4u} = \frac{1}{\sqrt{2}} \int_0^1 \frac{u^{-1/4}}{(1+u)^{3/2}} du.$$

Now make the substitution $v = 1/u$ in the integral $\int_1^\infty u^{-1/4}(1+u)^{-3/2} du$: $dv = -du/u^2$, $u^{-1/4} = v^{1/4}$, and $(1+u)^{-3/2} = v^{3/2}(1+v)^{-3/2}$, giving

$$\int_1^\infty u^{-1/4}(1+u)^{-3/2} du = \int_0^1 v^{-1/4}(1+v)^{-3/2} dv.$$

Therefore

$$\int_0^\infty u^{-1/4}(1+u)^{-3/2} du = 2 \int_0^1 u^{-1/4}(1+u)^{-3/2} du.$$

The left-hand side is the standard Beta-function integral

$$\int_0^\infty u^{\alpha-1}(1+u)^{-\beta} du = B(\alpha, \beta - \alpha) = \frac{\Gamma(\alpha)\Gamma(\beta - \alpha)}{\Gamma(\beta)}$$

with $\alpha = 3/4$, $\beta = 3/2$, giving $\Gamma(3/4)\Gamma(3/4)/\Gamma(3/2) = 2\Gamma(3/4)^2/\sqrt{\pi}$. Substituting back, we obtain

$$I_\infty = \frac{1}{\sqrt{2}} \cdot \frac{1}{2} \cdot \frac{2\Gamma(3/4)^2}{\sqrt{\pi}} = \frac{\Gamma(3/4)^2}{\sqrt{2\pi}}. \quad \square$$

5.2. Asymptotics of $s_0(a)$ and $r(a)$.

Lemma 5.2 (Leading-order asymptotics of $s_0(a)$). *As $a \rightarrow \infty$,*

$$(34) \quad s_0(a) = \log a + \log \frac{\sqrt{2}}{I_\infty} + o(1) = \log a + \log \frac{2\sqrt{\pi}}{\Gamma(3/4)^2} + o(1).$$

Proof. We analyze the free boundary equation (17):

$$\tanh \varphi(s_0) = \frac{B(s_0) K(a)}{a \sinh(2s_0)}.$$

For $a \rightarrow \infty$ and $s_0 \rightarrow \infty$, the right-hand side is

$$\frac{B(s_0) \cdot a}{a \sinh(2s_0)} (1 + O(a^{-2})) = \frac{\sqrt{a/2} e^{s_0} (1 + o(1))}{e^{2s_0}/2} (1 + o(1)) = \sqrt{2a} e^{-s_0} (1 + o(1)),$$

where we used $B(s)^2 \sim (a/2)e^{2s}$ and $\sinh(2s_0) \sim e^{2s_0}/2$ for large s_0 .

For the left-hand side, by (32), $\varphi(s_0) = I_\infty/\sqrt{a} \cdot (1 + o(1))$ provided $s_0 \rightarrow \infty$, hence $\varphi(s_0) \rightarrow 0$, and $\tanh \varphi(s_0) \sim \varphi(s_0) \sim I_\infty/\sqrt{a}$.

Equating the leading orders:

$$\frac{I_\infty}{\sqrt{a}} \sim \sqrt{2a} e^{-s_0} \iff e^{-s_0} \sim \frac{I_\infty}{a\sqrt{2}} \iff s_0 = \log \frac{a\sqrt{2}}{I_\infty} + o(1),$$

which is (34). □

Proof of Theorem 1.2. By Lemma 2.4,

$$\cosh r(a) = A(s_0) \cosh \varphi(s_0).$$

For $a \rightarrow \infty$ and $s_0 \rightarrow \infty$:

$$A(s_0)^2 \sim (a/2)e^{2s_0}, \quad \cosh \varphi(s_0) \rightarrow 1 \text{ (since } \varphi(s_0) \rightarrow 0).$$

Hence $\cosh r(a) \sim \sqrt{a/2} e^{s_0} (1 + o(1))$, so $r(a) = \log(2 \cosh r(a)) + o(1) = \log(\sqrt{2a} e^{s_0}) + o(1)$, that is,

$$r(a) = \frac{1}{2} \log(2a) + s_0 + o(1) = \frac{1}{2} \log 2 + \frac{1}{2} \log a + \log a + \log(\sqrt{2}/I_\infty) + o(1).$$

Combining the two terms:

$$r(a) = \frac{3}{2} \log a + \left(\frac{1}{2} \log 2 + \frac{1}{2} \log 2 - \log I_\infty\right) + o(1) = \frac{3}{2} \log a + \log 2 - \log I_\infty + o(1).$$

Substituting $I_\infty = \Gamma(3/4)^2/\sqrt{2\pi}$ from Lemma 5.1:

$$d_\infty := \log 2 - \log \frac{\Gamma(3/4)^2}{\sqrt{2\pi}} = \log \frac{2\sqrt{2\pi}}{\Gamma(3/4)^2}.$$

Using the reflection identity $\Gamma(1/4)\Gamma(3/4) = \pi\sqrt{2}$, this equals

$$d_\infty = \log \frac{\sqrt{2}\Gamma(1/4)^2}{\pi^{3/2}},$$

which is (6). The real-analyticity of $a \mapsto r(a)$ on $(1/2, \infty)$ follows from the implicit function theorem applied to (17), as the left-hand side is real-analytic in (s, a) and its s -derivative at $s = s_0(a)$ is non-zero. □

Proof of Proposition 1.3. Set $\epsilon := a - 1/2 > 0$ and look for $s_0(a)$ in the form $s_0 = \rho\sqrt{\epsilon}$, with $\rho = \rho(\epsilon)$ to be determined. The expansions

$$\cosh(2\rho\sqrt{\epsilon}) = 1 + 2\rho^2\epsilon + O(\epsilon^2), \quad \sinh(2\rho\sqrt{\epsilon}) = 2\rho\sqrt{\epsilon}(1 + O(\epsilon)),$$

together with $K(a) = \sqrt{\epsilon(1+\epsilon)} = \sqrt{\epsilon}(1 + O(\epsilon))$, give

$$B(s_0)^2 = (1/2 + \epsilon) \cosh(2s_0) - 1/2 = \epsilon(1 + \rho^2) + O(\epsilon^2), \quad A(s_0)^2 = 1 + \epsilon(1 + \rho^2) + O(\epsilon^2),$$

and consequently

$$(35) \quad R(s_0) = \frac{B(s_0)K(a)}{a \sinh(2s_0)} = \frac{\sqrt{\epsilon}\sqrt{1+\rho^2}\sqrt{\epsilon}}{(1/2+\epsilon) \cdot 2\rho\sqrt{\epsilon}}(1+O(\epsilon)) = \sqrt{\epsilon} \frac{\sqrt{1+\rho^2}}{\rho}(1+O(\epsilon)).$$

For the left-hand side of (17), the change of variable $t = \tau\sqrt{\epsilon}$ in the integral defining φ yields, using $A(t)^2 = 1 + O(\epsilon)$ and $B(t)^2 = \epsilon + t^2 + O(\epsilon^2)$ uniformly for $t = O(\sqrt{\epsilon})$,

$$(36) \quad \varphi(s_0) = \sqrt{\epsilon} \int_0^\rho \frac{d\tau}{\sqrt{1+\tau^2}}(1+O(\epsilon)) = \sqrt{\epsilon} \operatorname{arcsinh}(\rho)(1+O(\epsilon)),$$

hence $\tanh \varphi(s_0) = \sqrt{\epsilon} \operatorname{arcsinh}(\rho)(1+O(\epsilon))$. Equating this to (35) and dividing by $\sqrt{\epsilon}$,

$$(37) \quad \operatorname{arcsinh}(\rho) = \frac{\sqrt{1+\rho^2}}{\rho}(1+O(\epsilon)).$$

The function $F(\rho) := \operatorname{arcsinh}(\rho) - \sqrt{1+\rho^2}/\rho$ is real-analytic on $(0, \infty)$, strictly increasing (its first term is strictly increasing, the second is strictly decreasing), satisfies $F(\rho) \rightarrow -\infty$ as $\rho \rightarrow 0^+$ and $F(\rho) \rightarrow +\infty$ as $\rho \rightarrow \infty$. Therefore the limiting equation (8) has a unique positive root ρ_* , and by the implicit function theorem applied at ρ_* (where $F'(\rho_*) > 0$), $\rho(\epsilon) = \rho_* + O(\epsilon)$.

To translate into $r(a)$, recall $\cosh r(a) = A(s_0) \cosh \varphi(s_0)$. From the expansions above and (36),

$$A(s_0) = 1 + \frac{\epsilon}{2}(1 + \rho_*^2) + O(\epsilon^2), \quad \cosh \varphi(s_0) = 1 + \frac{\epsilon}{2} \operatorname{arcsinh}^2(\rho_*) + O(\epsilon^2).$$

At $\rho = \rho_*$ the relation (8) gives $\operatorname{arcsinh}^2(\rho_*) = (1 + \rho_*^2)/\rho_*^2$, and therefore

$$A(s_0) \cosh \varphi(s_0) = 1 + \frac{\epsilon}{2}(1 + \rho_*^2) \left(1 + \frac{1}{\rho_*^2}\right) + O(\epsilon^2) = 1 + \frac{\epsilon}{2}(\rho_* + \rho_*^{-1})^2 + O(\epsilon^2).$$

Using $\operatorname{arccosh}(1+x) = \sqrt{2x}(1+O(x))$ as $x \rightarrow 0^+$,

$$r(a) = \operatorname{arccosh}(A(s_0) \cosh \varphi(s_0)) = (\rho_* + \rho_*^{-1})\sqrt{\epsilon}(1+O(\epsilon)),$$

which is (7). □

Remark 5.3. We do not give a closed form for $r'(a)$ in terms of elementary functions for general a . An analytic monotonicity proof of $a \mapsto r(a)$ on $(1/2, \infty)$ can in principle be obtained from the implicit function theorem applied to (17) by checking the sign of $\partial s_0/\partial a$, but is not pursued here.

6. OPEN PROBLEMS

The analysis of this note suggests several natural directions for further investigation.

6.1. Mode-by-mode index decomposition. Corollary 4.4 establishes that the contribution of mode $|k| = 1$ to the Robin Morse index of Σ_a equals exactly 2. Combined with the lower bound $\operatorname{ind}(\Sigma_a) \geq 4$ of Medvedev [8], the contribution of modes $|k| \neq 1$ to the index is ≥ 2 . The conjectural equality $\operatorname{ind}(\Sigma_a) = 4$ of Medvedev would follow from a sharper mode decomposition, which we formulate explicitly:

Conjecture 6.1 (Mode-by-mode Morse index of Σ_a). *For every $a > 1/2$:*

- (a) *the Robin Morse index of Σ_a in mode $|k| = 0$ equals 2;*
- (b) *the Robin Morse index of Σ_a in mode $|k| \geq 2$ equals 0.*

In particular, combined with Corollary 4.4, this would give $\operatorname{ind}(\Sigma_a) = 4$ and $\operatorname{nul}(\Sigma_a) = \operatorname{nul}(\Sigma_a)|_{|k|=1} = 2$ for every $a > 1/2$, settling Medvedev's conjecture.

The conjecture is supported by two structural facts:

- (i) In the Euclidean critical catenoid (Devyver [2]), the analogous mode decomposition holds: mode $|k| = 0$ contributes 2 to the index, mode $|k| = 1$ contributes 2, and modes $|k| \geq 2$ contribute 0.
- (ii) The radial mode- k Jacobi operator has potential $V_k(s) = |\mathbb{II}|^2 - 2 - k^2/B(s)^2 = 2K^2/B(s)^4 - 2 - k^2/B(s)^2$ which becomes strictly more negative (hence more stabilizing on the μ -spectrum side) as $|k|$ grows.

The proof of (a) appears delicate, since unlike mode $|k| = 1$ no Killing–Jacobi field exists in mode $|k| = 0$ (the Killing fields L_{12}, L_{13}, L_{23} all lie in modes $|k| \leq 1$). A natural candidate for an explicit second-order test function in mode $|k| = 0$, by analogy with (23), would be ∂_a -derivatives of Φ_a^A along the family $\{\Sigma_a\}_a$; making this rigorous, together with a Sturm-theoretic count, is left to future work.

6.2. Higher-order asymptotics of the boundary radius. Theorem 1.2 gives $r(a) = \frac{3}{2} \log a + d_\infty + o(1)$, and Remark 1.6 (in the proof) shows that the remainder is in fact $O(a^{-1})$. We pose:

Question 6.2. *Compute, in closed form involving values of Γ at rational arguments, the next coefficient d_1 in the expansion*

$$r(a) = \frac{3}{2} \log a + d_\infty + \frac{d_1}{a} + o(a^{-1}) \quad (a \rightarrow \infty).$$

By the same Beta-function machinery used in §5, d_1 should be expressible as a definite integral of the form $\int_0^\infty \cosh(2t)^{-5/2} dt$ or a linear combination of such integrals, all of which evaluate in closed form via $B(p, q) = \int_0^\infty \cosh(2t)^{-(p+q)} (\sinh 2t)^{2q-1} dt$ at rational p, q .

6.3. Strict monotonicity of $r(a)$. Theorem 1.2 establishes real-analyticity of $a \mapsto r(a)$ on $(1/2, \infty)$. Numerical investigation strongly suggests:

Question 6.3. *Is $a \mapsto r(a)$ strictly increasing on $(1/2, \infty)$?*

By the implicit function theorem applied to (17), this reduces to showing that $\partial F/\partial a > 0$ at the implicit solution, where $F(a, s_0)$ denotes the free boundary equation. The sign of this partial derivative does not seem to follow from elementary manipulations.

6.4. Higher-dimensional analogs. The Mori family of catenoids extends to \mathbb{H}^{n+1} for $n \geq 2$ (cf. do Carmo–Dajczer [3]), each with a critical free boundary representative $\Sigma_a^{(n)} \subset B^{n+1}(r^{(n)}(a))$.

Question 6.4. *Do the closed forms of Lemma 3.4, Corollary 3.5 and Theorem 1.2 extend to \mathbb{H}^{n+1} ? Specifically:*

- (a) *Does the Killing–Jacobi profile in the rotation modes admit a closed form $\partial_s \Phi^A$ for some Minkowski coordinate A ?*
- (b) *Does the asymptotic constant $d_\infty^{(n)}$ admit a closed form involving values of Γ at rational arguments?*

For (b), heuristics suggest $d_\infty^{(n)} \propto \log(\Gamma(\alpha_n))$ for some explicit rational α_n depending on n , by the same Beta-function evaluation; the natural candidate is $\alpha_n = (2n - 1)/(4n - 2)$ or a closely related fraction.

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